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Atomic and electronic structures of amorphous $ZrO₂$ and $HfO₂$ films

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Abstract

Atomic and electronic structures of $ZrO₂$ and $HfO₂$ films deposited using sputtering technique were studied by X-ray diffraction (XRD), EXAFS spectroscopy, photoemission (with X-ray and ultraviolet excitation sources) and electron energy loss spectroscopy (EELS). XRD results indicated that the as-deposited films were amorphous and a monoclinic phase was detected after annealing them in ambient at 800 °C. Photoemission results indicated that the $ZrO₂$ valence band consists of two bands separated by an ionic gap of 6 eV. The lower band of 10 eV width is occupied mainly by O 2s states. EELS results indicated that the bandgap of $ZrO₂$ and HfO₂ films are 4.7 and 5.7 eV, respectively. $© 2005$ Published by Elsevier B.V.

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1. Introduction

The continuous shrinking of dimensions of complementary metal-oxide-silicon (CMOS) devices has led to a need to replace the silicon dioxide

Corresponding author. Fax: +852 2788 7791. E-mail address: heiwong@ieee.org (H. Wong). by higher dielectric constant insulators. Because of high dielectric constant and potential for forming device grade interfaces with Si substrate, hafnium oxide (HfO₂) and zirconium oxide (ZrO₂), with similar chemistries, have emerged as promising replacements for $SiO₂$ [\[1,2\]](#page-5-0). However, when compared with the conventional silicon oxide, high- k interfaces appear to have high trap density, large

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leakage current and low surface mobility [\[1,3\]](#page-5-0). Since carrier transport properties such as carrier injection, trapping, trap ionization or charge transport in the dielectric bulk are governed by atomic and electronic structure of the material, it is crucial to understand these important features. In this work, we explore various atomic and electronic structures of sputtered and annealed $ZrO₂$ and HfO_2 films by using several analytical tools.

2. Experimental

The zirconium and hafnium oxide was deposited on (100) *n*-type Si substrate with resistivity of 20Ω cm using an ARC-12M sputtering machine. Before the sputtering, a standard RCA cleaning process was conducted to clean the sample. Zirconium and hafnium metal of 99% purity were used as targets for sputtering. The argon and oxygen gases were introduced into the chamber after it was pumped to 10^{-6} torr. The argonto-oxygen flow ratio was fixed at 2:23 and the pressure of the chamber was kept at 10^{-3} torr during the sputtering. The X-ray diffraction (XRD) measurements were performed with a high-resolution powder diffractometer available in Budker Institute of Nuclear Physics in Novosibirsk. Zirconium K edge (Zr K) EXAFS spectra were measured using the synchrotron radiation (SR) of the VEPP-3 storage ring at the same Institute. The photoemission studies were performed with a VG ESCALAB 220i-XL photoelectron spectroscopy system with a monochromatic aluminum $K\alpha$ source (1486.6 eV). An unfiltered He I (21.2 eV) gas-discharge lamp was used for ultraviolet photoelectron spectroscopy (UPS) study. Electron energy loss spectra (EELS) were measured using LAS-3000 Auger Electron Spectrometer.

3. Results and discussion

Fig. 1 shows the XRD spectra of both an asdeposited $HfO₂$ film and a sample subjected to 800 °C annealing in atmosphere. The as-deposited film shows an amorphous structure, whereas the thermal annealing results in the appearance of

Fig. 1. X-ray diffraction (XRD) spectra of $HfO₂$ films: (1) as deposited, (2) annealed at 800 $^{\circ}$ C for 30 min.

the tetragonal phase. EXAFS measurement was used to study the structure peculiarities. Fig. 2 shows the radial distribution function of $HfO₂$. Both virgin and annealed films show pronounced peaks in the first coordination sphere corresponding to the Hf–O bond. The peak intensity increases after the thermal annealing (see trace 2). No pronounced Hf–Hf peak is found in virgin sample but in annealed $HfO₂$, the Hf–Hf peak is enhanced significantly. These results indicate that the asdeposited film is amorphous and it changes to

Fig. 2. Radial distribution function (normalized Zr K X-ray absorption coefficient) of as-deposited $ZrO₂$ film (trace 1) and of the sample annealed thermally at 800 $^{\circ}$ C (trace 2).

monoclinic phase after the 800 $^{\circ}$ C annealing. Fig. 3 shows the EXAFS of as sputtered and annealed ZrO₂ films. Similar phenomena are observed. The thermal annealing results in the increase of Zr–O peak intensity so that the Zr–Zr peak in the annealed sample becomes comparable to the Zr–O peak.

Band-structure program BAND which is based on Amsterdam Density Functional (ADF) was used to calculate the electronic structure of cubic $ZrO₂$. The spin restricted option of the program was used and no relativistic effects were taken into account. The exchange-correlation potential was described by the local density approximation (LDA). Combination of numerical atomic orbitals (NOs) from Hermann–Skillman type free atom calculation and Slater type atomic orbitals (STO) Zr 4p, 4d, 5s and O 2s, 2p, 53 functions per unit cell were treated as the valence states. Atomic NOs of core states were considered frozen. Other details of the ADF simulation can be found in the work by Velde and Baerends [\[4\].](#page-5-0)

The lower valence band of $ZrO₂$ is mainly made up with oxygen 2p states and only in the lower energy part of the top band the Zr 4d states give remarkable contribution. The top of upper valence

Fig. 3. Radial distribution function (normalized Hf K X-ray absorption coefficient) of as-deposited $HfO₂$ film (trace 1) and of the sample annealed thermally at 800 \degree C (trace 2).

band located at X point of Brillouin zone. The lowest conduction band consists mainly the Zr 4d states. The bottom of the band is located at Γ point of the Brillouin zone. From the dispersion law, nearby Γ point along ΓX -, ΓL -, and ΓK directions, the electron effective masses were obtained. In Γ , point there are two generated states which are split along ΓX - and ΓK -directions, and remain generated along ΓL -direction. Corresponding to these directions, we have five different electron masses with m_e^*/m_0 value in the range of 0.73–2.26. The hole masses at X point are in the range of 0.33–3.33; namely, both light and heavy holes existed.

The valence band electronic structure of the amorphous $ZrO₂$ and $HfO₂$ were studied using XPS and UPS. According to the experimental data and the quantum-chemical simulation, the $ZrO₂$ valance band consists of two bands: the upper band mainly consists of O 2p states with mixture of Zr 4d, Zr 5s states, and the lower band with O 2s states (see Fig. 4). Similar $ZrO₂$ valence band electronic structure was reported by Foster et al. [\[5,6\].](#page-5-0) The ionization cross-section of electron in the valence band depends on their wave function symmetry and the excitation energy [\[7\]](#page-5-0). At the excitation energy of 1486.6 eV, the main contribution to the electron ionization cross-section in the top valence band is due to Zr 4p and Zr 4d

Fig. 4. Experimental X-ray photoelectron spectra (solid line) of amorphous $ZrO₂$ valence band (solid line) and calculated XPS valance band spectra of cubic zirconia (dashed line).

electrons. For excitation energy of 40.8 eV, the largest contributions to the electron ionization are from O 2p and Zr 4p electrons; and at excitation energy of 21.2 eV, the ionization cross-section is contributed mainly from O 2p and Zr 4d electrons.

Ultraviolet photoelectron spectrum of $ZrO₂$ measured at energy 21.2 eV is shown in Fig. 5 together with a gold spectrum for calibration. The low-energy threshold of UPS spectra is due to the high-energy photoelectrons. The low energy tail of UPS spectra contains information of the top position of $ZrO₂$ valence band. Since the work function of gold is 5.1 eV and the top of $ZrO₂$ valence band is 2.5 eV above the gold Fermi level, it follows that the position of the amorphous $ZrO₂$ valence band relative to the vacuum level is 7.6 eV. The gold Fermi level is same as the top of silicon valence band, the valence band offset (hole barrier) at $Si/ZrO₂$ interface is thus 2.5 eV.

This value agrees with the hole barrier determined using internal photoemission and photoconductivity experiments for atomic layer deposited (ALD) $ZrO₂$ [\[8,9\].](#page-5-0)

The energy gap of $HfO₂$ and $ZrO₂$ were investigated with electron energy loss spectroscopy (EELS) measurement at electron beam energy of 200 eV and the results are shown in Fig. 6. In these diagrams, the zero energy corresponds to the elastic electron peak and negative energies correspond to the inelastic energy losses. For amorphous $ZrO₂$ prepared using sputtering method, the measured bandgap is 4.7 eV (see Fig. $6(a)$) which is in the

Fig. 5. Ultraviolet photoelectron spectroscopy of amorphous $ZrO₂$ and gold valence band excited by He I line (21.2 eV).

Fig. 6. Electron energy loss spectra of: (a) amorphous $ZrO₂$ and (b) amorphous $HfO₂$ at electron beam energy 200 eV. Zero energy corresponds to the elastic electron peak and negative energies correspond to the inelastic energy losses.

mostly reported range [\[9–11\]](#page-5-0) and differs from the 5.8 eV value given by Robertson [\[12\]](#page-5-0). For hafnium oxide, the measured bandgap is 5.7 eV (see [Fig.](#page-3-0) $6(b)$ $6(b)$). This value is in the range reported by a recent study using photoelectron emission measurements [\[13\]](#page-5-0) and slightly less than the 6.0 eV value reported earlier [\[12\]](#page-5-0). Yet the differences may be due to the oversimplified procedure for extracting the bandgap data. A better approach by subtracting the zero-loss peak and restoring the loss function to determine the onset of band-to-band excitations [\[14\]](#page-5-0) is under investigation.

Current–voltage characteristics of the p-Si/ $ZrO₂/Al$ structure measured in dark and at illumination are shown in Fig. 7. At positive biasing, current saturation due to minority carrier (electron) injection from silicon into dielectric is observed. This result indicates that the current conduction is mainly contributed by electron for positive biasing on Al electrode. In the n-Si/ $ZrO₂/Al$ structure (see Fig. 8), the current saturation occurs at negative biasing on the Al electrode and suggests that the current conduction is contributed by hole current. Hence, the current con-
duction in $ZrO₂$ is two bands resembling to the

Fig. 7. Current–voltage characteristics of p-Si $ZrO₂/Al$ structure at both polarities in dark (solid line) and under illumination (dashed curve). The current saturation at positive Al biasing indicates the electron injection from Si into the $ZrO₂$ film.

Fig. 8. Current–voltage characteristics of n-Si/ZrO₂/Al structure at both polarities in dark in under illumination. The current saturation at negative Al biasing indicates the hole injection from Si into $ZrO₂$.

silicon nitride case.

4. Conclusion

In summary, the atomic and electronic structures of sputtered $ZrO₂$ and $HfO₂$ were investigated. Thermal annealing at 800 °C results in transferring the amorphous phase to tetragonal ones and narrow Zr–O and Hf–O EXAFS peaks were observed. Photoemission measurements indicated that the $ZrO₂$ and $HfO₂$ valence band consists of two bands separated by an ionic gap of 6 eV. The bandgaps of $ZrO₂$ and $HfO₂$ estimated from EELS are 4.7 and 5.7 eV, respectively.

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