## **Oxidation states in scanning-probe-induced Si3N4 to SiO<sup>x</sup> conversion studied by scanning photoemission microscopy**

R. Klauser<sup>a)</sup> and I.-H. Hong

*Synchrotron Radiation Research Center, No. 1 R&D Road VI, Hsinchu 300, Taiwan, Republic of China*

H.-J. Su, T. T. Chen, and S. Gwo

*Department of Physics, National Tsing-Hua University, Hsinchu 300, Taiwan, Republic of China*

S.-C. Wang

*Institute of Atomic and Molecular Sciences, Academia Sinica, Taipei 106, Taiwan, Republic of China*

T. J. Chuang<sup>b)</sup>

*Center for Condensed Matter Sciences, National Taiwan University, Taipei 106, Taiwan, Republic of China*

## V. A. Gritsenko

*Institute of Semiconductor Physics, Lavrentieva Avenue 13, Novosibirsk 630090, Russia*

(Received 29 June 2001; accepted for publication 22 August 2001)

The biased conductive probe of an atomic force microscope can induce local oxidation in ambience for converting silicon nitride films to silicon oxides with high reaction rate. Spatially resolved photoemission analysis with submicron resolution has been utilized to study the oxidation states of converted silicon oxide patterns in comparison with the surrounding  $Si<sub>3</sub>N<sub>4</sub>$  layer. The core level shift of the  $Si 2p$  photoelectron peak and the spectral features in the valence band reveal a complete conversion of silicon nitride to silicon oxide at a bias voltage of 10 V, with no remaining nitrogen left. The major oxide is  $SiO_2$ . The observed oxidation states of  $Si^{4+}$ ,  $Si^{3+}$ , and  $Si^{2+}$  show a gradient depth distribution indicating excess silicon in the layer. © *2001 American Institute of Physics.* [DOI: 10.1063/1.1415415]

The need for lithographic techniques with nanometer spatial resolution is obvious for the development of new quantum devices and molecular electronics. In the recent decade, scanned probe microscopy, such as scanning tunneling microscopy (STM) or atomic force microscopy (AFM) has emerged as an important tool in nanoscale device fabrication by utilizing the effect that a local high electrical field between the probing tip and sample can modify the surface chemistry in a selective way with nanometer scale resolution. This has been demonstrated for the AFM-induced anodic oxidation in air of  $Si$ ,<sup>1,2</sup> metal surfaces<sup>3</sup> and insulating thin films,4 as well as for the hydrogen-terminated Si surfaces by ultrahigh vacuum (UHV) STM-induced hydrogen bond breaking, and subsequent reaction with various molecules.<sup>5,6</sup> The recently reported nanoscale conversion of silicon nitride film to silicon oxide by one of the authors  $(S.G.)$  showed an amazingly high reaction rate, $7$  compared with the case of thermal oxidation. $8$  It is widely accepted that the mechanism for the process of Si anodic AFM oxidation in regular atmosphere at room temperature is essentially controlled by air humidity, i.e., the adsorption of water on the surface. Negatively charged oxygen ions are injected from the water, and under the high electric field, oxyanions are transported through the oxide layer to the  $Si/SiO<sub>2</sub>$  interface, where they react with the positively charged holes and Si to form  $SiO<sub>2</sub>$ . The oxide density is shown to be sensitive to voltage pulse

parameters and substrate doping.<sup>9</sup> Space charge effects associated with defect generation during growth are responsible for nonuniformity of the oxide. In the case of  $Si<sub>3</sub>N<sub>4</sub>$  conversion, the reaction between oxyanions, silicon nitride, and the holes accumulated at the  $Si<sub>3</sub>N<sub>4</sub>/oxide$  interface causes a replacement of nitrogen by oxygen. How far and under what condition  $Si_3N_4$  film can be completely converted to  $SiO_2$  or only partially into  $SiO_{x \leq 2}$  or  $SiO_xN_y$  is the subject of this letter and essential for understanding the silicon nitride oxidation mechanism. The conversion to  $SiO<sub>x<2</sub>$  would imply the coexistence of Si–O and Si–Si bonds in the oxide layer due to an excess of Si. Previous micro-Auger electron spectroscopy measurements of AFM-induced oxide pads on Si3N4 film showed the disappearance of the N-*KLL* peak and an enhancement of the O-KLL signal.<sup>4,7</sup> This indicates that the conversion to oxynitrides with a substantial amount of nitrogen in the layer is insignificant. Auger electron spectroscopy in this case, however, can only reveal elemental composition but cannot distinguish between chemical states. Photoelectron spectroscopy is known for chemical sensitivity, but generally has a lack of spatial resolution in the conventional mode of operation. We have utilized one of the few available systems in the world to perform spatially resolved photoemission experiments<sup>10</sup> with submicron resolution to study pads of AFM-converted silicon oxide patterned on  $Si<sub>3</sub>N<sub>4</sub>$  film.

The silicon nitride films were grown on  $p$ -type  $Si(001)$ wafers with the low pressure chemical vapor deposition technique.<sup>7</sup> The original film thickness was 50 Å, which was then reduced to about 25 Å by HF  $(1%)$  etching to guarantee the complete oxidation of the  $Si<sub>3</sub>N<sub>4</sub>$  film. Local field-induced

**Downloaded 31 Oct 2001 to 210.241.28.12. Redistribution subject to AIP license or copyright, see http://ojps.aip.org/aplo/aplcr.jsp**

a)Author to whom correspondence should be addressed; electronic mail: klauser@srrc.gov.tw

b) Also with: Institute of Atomic and Molecular Sciences, Academia Sinica, Taipei 106, Taiwan, Republic of China.

<sup>0003-6951/2001/79(19)/3143/3/\$18.00 © 2001</sup> American Institute of Physics 3143



FIG. 1. 16 simultaneously acquired scanning photoemission microscopy images of Si 2*p* photoelectrons at different kinetic energies from an 8  $\mu$ m×10  $\mu$ m pad of AFM-oxidized Si<sub>3</sub>N<sub>4</sub> film (image size: 21  $\mu$ m×21  $\mu$ m, photon energy: 384 eV).

oxidation was performed in air atmosphere with a commercial AFM system using conductive PtIr-coated probes. The applied bias voltages were between 7 and 10 V, producing various patterns of stripes and pads. For locating the pattern on the sample, however, large squares  $8 \mu m \times 10 \mu m$  were often used in this study. Without any further surface treatment, the samples were imaged in a UHV chamber by soft x-ray scanning photoemission microscopy (SPEM). This custom-made apparatus at the Synchrotron Radiation Research Center in Hsinchu, Taiwan is based on focusing a soft x-ray beam in the range between 250 and 800 eV from an undulator beamline by means of Fresnel zone plate optics onto a spot 0.1  $\mu$ m in size.<sup>11</sup> Details of the SPEM design and performance were given previously.<sup>10,11</sup> The overall spatial resolution in this experiment is 0.2  $\mu$ m.

Figure 1 shows a set of 16 simultaneously acquired photoemission images for a photon energy (PE) of 384 eV and the photoelectron binding energy (BE) between 96 and 108 eV, where the Si 2*p* peak is located. The pad size of AFMoxidized  $Si_3N_4$  film (bias voltage of 10 V) is 8  $\mu$ m×10  $\mu$ m. The gray scale reflects the intensity of the photoelectrons. Obviously, the oxidized pattern can be clearly distinguished. The  $Si 2p$  signal inside the pad is maximal at channel 7, representing the state of silicon oxide. Channel 9, with 1.5 eV higher kinetic energy than channel 7, shows reversed contrast and reflects the second chemical state of Si, i.e.,  $Si<sub>3</sub>N<sub>4</sub>$ . Channels of background intensity far beyond the peak signal (channels  $1-4$  and  $13-16$ ) reveal very little contrast. The topographic height difference between the oxide pattern and  $Si<sub>3</sub>N<sub>4</sub>$  is too small (the protruded oxide is about  $60\%$  of the  $Si<sub>3</sub>N<sub>4</sub>$  film thickness<sup>4</sup>) to become evident in the SPEM image. Similar SPEM maps as for Si 2*p* can be obtained for the energy ranges of N 1*s* and O 1*s* core levels.



FIG. 2. Microphotoemission spectra (dotted curves) and peak analysis of  $Si 2p$  on the  $Si<sub>3</sub>N<sub>4</sub>$  film and inside the AFM-oxidized pad. Insets are the corresponding N 1*s* spectra.

From the imaging mode, one can switch to the spectroscopic mode for micro-ultraviolet or micro-x-ray photoelectron spectroscopy analysis by selecting a particular point on the image and measuring photoemission energy distribution curves. This is displayed in Fig. 2 for a point near the center of the oxidized pad and for two different photon energies. The photon energy has been calibrated by the Fermi edge of a gold sample attached to the same sample holder. Avoiding the possible calibration error due to charging, the binding energy scale has also been determined by the carbon contamination of the samples. After subtraction of the secondary electron background, the Si 2*p* spectra were fitted by doublets of Voigt function with a standard spin-orbit splitting and branching ratio of the Si  $2p_{3/2}$  and Si  $2p_{1/2}$  peak constituents of 0.61 eV and 2, respectively. The Si 2*p* spectra of the silicon oxide pad for both photon energies show a clear asymmetry to lower BE. The three deconvoluted components have Si  $2p_{3/2}$  peak maxima at 103.3, 102.4, and 101.6 eV with full-width-at-half-maximum  $(FWHM)$  of 1.3, 1.1, and 1.1 eV, respectively. The BE and the width of the major component agree well with reference values found in the literature for  $SiO_2$ .<sup>12</sup> Si 2p core-level spectroscopy of thermally grown  $SiO<sub>2</sub>$  film on  $Si(100)$  and  $Si(111)$  surfaces exhibits four oxidation states of silicon. Besides  $SiO_2(Si^{4+})$ , intermediate states of  $Si^{1+}(Si_2O)$ ,  $Si^{2+}(SiO)$ , and  $Si^{3+}(Si_2O_3)$  exist at the SiO<sub>2</sub>/Si interface. The energy positions of the two deconvoluted shoulder peaks of the SiO*<sup>x</sup>* square in Fig. 2 relative to the  $SiO<sub>2</sub>$  peak agree with the

**Downloaded 31 Oct 2001 to 210.241.28.12. Redistribution subject to AIP license or copyright, see http://ojps.aip.org/aplo/aplcr.jsp**



FIG. 3. Valence band spectra on the  $Si<sub>3</sub>N<sub>4</sub>$  film, inside the AFM-oxidized pad, and of a 60 Å thick  $SiO<sub>2</sub>$  reference sample (thermally oxidized). Energy zero point corresponds to the Fermi edge of an Au reference sample.

suboxides  $Si^{3+}$  and  $Si^{2+}$  for the thermally oxidized Si film, but with about 20% larger peak widths. The  $Si^{3+}$  species increases in intensity for larger probing depth at photon energy of 622 eV (the electron escape length at  $PE=384$  eV is about 10 Å and at PE=622 eV is about 14 Å), whereas the  $Si<sup>2+</sup>$  suboxide intensity remains small. We conclude that those chemical states are due to intermediate oxides produced by the AFM-induced conversion of  $Si<sub>3</sub>N<sub>4</sub>$ , but with no remaining nitrogen incorporated in the film as indicated by the absence of N 1*s* (insert in Fig. 2) and N 2*s* (Fig. 3) signals. The silicon oxide pad is by about 1.5 nm thicker than the nitride film because of the additional protrusion. The suboxides can therefore not directly originate from the  $SiO<sub>2</sub>/Si$ interface, but from an extended area in the layer containing excess Si–Si bonds.

The Si  $2p$  peak of  $Si<sub>3</sub>N<sub>4</sub>$  film is quite broad with contributions from stoichiometric  $Si<sub>3</sub>N<sub>4</sub>$  with  $Si<sub>2</sub>p<sub>3/2</sub>$  peak maximum at 101.6 eV and from native oxide contamination at the surface with the corresponding peak at 102.6 eV. For PE =622 eV, additional contributions from Si at the  $Si_3N_4/Si$ interface with 100.7 eV BE is also evident. The difference between  $SiO<sub>2</sub>$  and the oxygen contamination of  $Si<sub>3</sub>N<sub>4</sub>$  film also appears in the binding energy values of O 1*s* with 533.2 eV for silicon dioxide and 532.8 eV for native oxide. Both the N 1*s* signal with maximum intensity at 398.1 eV BE and the fitted Si  $2p_{3/2}$  peak at 101.6 eV have FWHM of 1.4 eV, as expected for a well-prepared  $Si<sub>3</sub>N<sub>4</sub>$  film. Figure 3 shows the valence band spectra of a silicon nitride film, the AFMoxidized pad, and a 60 Å thick reference sample of thermal  $SiO<sub>2</sub>$  for two different photon energies. The main contributions in these spectra between 15 and 35 eV are from O 2*s* and N 2*s*, and those below 15 eV are mostly from Si 3*s*, Si  $3p$ ,  $02p$ , and  $N2p$ . The AFM-induced oxidation results in a strong increase in O 2*s* intensity and the absence of N 2*s* signal. The band edge and the shape of the valence band resemble the spectrum of thermal dioxide. However, a slight increase in background intensity between 1 and 4 eV illustrates the presence of excess silicon in the layer. The band offset between 60 Å  $SiO_2$  and  $Si<sub>3</sub>N<sub>4</sub>$  film is about 4.0 eV. This value can be compared with 1.5 eV, measured by hole injection in  $Si<sub>3</sub>N<sub>4</sub>$ .<sup>13</sup> We can assume that this inconsistency can be related to the existence of the silicon-rich oxynitride phase at the  $Si_3N_4$ /native oxide interface, which was previously observed.<sup>8</sup>

In summary, spatially resolved soft x-ray photoemission microscopy has been employed to study AFM-induced oxidation on silicon nitride films. This technique provides direct evidence of the oxidation state distribution. It is found that a well-prepared sample can be completely converted to silicon oxide. The major oxide is  $SiO<sub>2</sub>$ . Within the probing depth of the photoelectrons, other oxidation states of  $Si^{3+}$  and  $Si^{2+}$ are also observed. Apparently, suboxides containing Si–Si bonds are produced in the layer with gradient depth distribution. Excess Si in the layer escaping from the  $Si_3N_4/\text{o}x$  ide interface might explain the enhancement of the oxidation reaction.

The authors acknowledge Dr. T.-S. Chao of National Nano Device Laboratory for supplying the  $Si<sub>3</sub>N<sub>4</sub>$  and  $SiO<sub>2</sub>$ samples. This work is supported by the National Science Council of ROC and the Ministry of Education.

- <sup>1</sup> J. A. Dagata, J. Schneir, H. H. Harary, C. J. Evans, M. T. Postek, and J. Bennett, Appl. Phys. Lett. **56**, 2001 (1990).
- 2M. Calleja, J. Anguita, R. Garcia, K. Birklund, F. Perez-Murano, and J. A. Dagata, Nanotechnology **10**, 34 (1999).
- $3$ E. S. Snow and P. M. Campbell, Science  $270$ , 1639 (1995).
- 4F. S.-S. Chien, J.-W. Chang, S.-W. Lin, Y.-C. Chou, T. T. Chen, S. Gwo, T.-S. Chao, and W.-F. Hsieh, Appl. Phys. Lett. **76**, 360 (2000).
- 5T.-C. Shen, C. Wang, J. W. Lyding, and J. R. Tucker, Appl. Phys. Lett. **66**, 976 (1995).
- 6M. C. Hersam, N. P. Guisinger, and J. W. Lyding, Nanotechnology **11**, 70  $(2000).$
- <sup>7</sup>F. S.-S. Chien, Y.-C. Chou, T. T. Chen, W.-F. Hsieh, T.-S. Chao, and S. Gwo, J. Appl. Phys. **89**, 2465 (2001).
- 8V. A. Gritsenko, S. N. Svitasheva, I. P. Petrenko, H. Wong, J. B. Xu, and I. H. Wilson, J. Electrochem. Soc. **146**, 780 (1999).
- 9K. Morimoto, F. Perez-Murano, and J. A. Dagata, Appl. Surf. Sci. **158**,  $205$   $(2000)$ .
- 10C.-H. Ko, R. Klauser, D.-H. Wei, H.-H. Chan, and T. J. Chuang, J. Synchrotron Radiat. **5**, 299 (1998).
- <sup>11</sup> I.-H. Hong, T.-H. Lee, G.-C. Yin, D.-H. Wei, J.-M. Juang, T.-E. Dann, R. Klauser, T. J. Chuang, C. T. Chen, and K.-L Tsang, Nucl. Instrum. Methods Phys. Res. A 905, 467 (2001).
- 12F. J. Himpsel, F. R. McFeely, A. Taleb-Ibrahimi, J. A. Yarmoff, and G. Hollinger, Phys. Rev. B 38, 6084 (1988).
- 13V. A. Gritsenko, E. E. Meerson, and Yu. N. Morokov, Phys. Rev. B **57**, R2081 (1998).